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# The influence of the Coulomb explosion on the energy loss of $H_2^+$ and $H_3^+$ molecules channeling along the Si $\langle 100 \rangle$ direction

R.C. Fadanelli <sup>a,\*</sup>, P.L. Grande <sup>a</sup>, M. Behar <sup>a</sup>, J.F. Dias <sup>a</sup>, G. Schiwietz <sup>b</sup>

<sup>a</sup> Instituto de Física da Universidade Federal do Rio Grande do Sul, Avenida Bento Gonçalves 9500, C.P. 15051, CEP 91501–970, Porto Alegre, Brazil

<sup>b</sup> Hahn-Meitner-Institut, Abteilung SF4, Glienicker Strasse 100, D-14109 Berlin, Germany

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## Abstract

In this work we have measured the contribution of the Coulomb explosion to the electronic stopping power of molecular hydrogen ions ( $H_2^+$  and  $H_3^+$ ) channeling along the Si  $\langle 100 \rangle$  direction. We have used a SIMOX target, consisting of crystalline Si  $\langle 100 \rangle$  with a buried layer of  $SiO_2$ . The measurements of the energy loss of  $H^+$ ,  $H_2^+$  and  $H_3^+$  have been carried out using the standard channeling Rutherford backscattering spectrometry. The energy loss has been measured around the Si  $\langle 100 \rangle$  channel at a fixed energy per nucleon (150 keV/amu) as a function of the tilt and azimuthal angles. The present results show the effect of Coulomb explosion, which enlarges the protons traversal energy and consequently the channeling energy loss. This heating effect due to  $H_3^+$  ions is about two times larger than  $H_2^+$  molecules and amounts to about 5% of the total stopping power.

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## 1. Introduction

Beams of ionic clusters are useful tools in both fundamental and applied researches, namely mate-

rial science, plasma and nuclear physics. It is well established that the effects of a molecular beam clearly deviates from those related with its individual components. The first evidence of interference effects among the projectile constituents, namely the vicinage effect, was first reported by Brandt et al. [1] and it was further studied by several authors, e.g. [2,3]. A detailed review of the subject is given in [4].

\* Corresponding author. Tel.: +55 51 3316 6517; fax: +55 51 3316 7286.

E-mail address: [raul@if.ufrgs.br](mailto:raul@if.ufrgs.br) (R.C. Fadanelli).

A swift molecular beam, when entering a solid, shows a second effect. The projectile loses its bonding electrons in the target and its components undergo a molecular breakup process due to quasi-Coulomb repulsive forces. This is the so-called Coulomb explosion (see, for example, [5] for experiments and [6] for its theoretical description).

When a molecular beam enters a crystal under channeling conditions, these two effects may compete. From one hand, the vicinage effect leads to a non-additive stopping power, i.e. an enhancement or a reduction dependent on the projectile speed. On the other hand, the Coulomb explosion of the molecule tends to enlarge the transversal energy of the components (the so called transverse Coulomb heating [7,8]), increasing consequently the total stopping power. The transverse Coulomb heating can occur only under channeling conditions, where it affects the ion flux distribution. In principle, both effects cannot be separated and therefore, the interplay between them in the stopping power remains still unclear. In fact, differences in stopping power results of  $H^+$  ions and  $H_2^+$  and  $H_3^+$  molecules have been attributed alternatively to vicinage effects [4] or to the Coulomb explosion of the molecules [9].

In order to study the role of the Coulomb explosion in the energy-loss processes, we have undertaken the present experiment where, by measuring the electronic stopping power of  $H^+$ ,  $H_2^+$  and  $H_3^+$  under channeling conditions as a function of the tilting angle, we were able not only to isolate the effect of the Coulomb explosion but also to estimate its contribution to the total stopping power.

## 2. Experimental setup and procedure

The measurements of the  $H^+$ ,  $H_2^+$  and  $H_3^+$  stopping powers were done using a SIMOX target composed by a 200 nm Si  $\langle 100 \rangle$  film on top of a 400 nm  $SiO_2$  film, both being constructed over a Si  $\langle 100 \rangle$  wafer. Beams of  $H^+$ ,  $H_2^+$  and  $H_3^+$  of 150 keV/amu incident on the SIMOX target were backscattered in the target and collected by a surface barrier detector. The resolution of the detector plus the electronic system was better than 7 keV.

In the first place a  $H^+$  channeling spectrum was obtained. With this aim we first located the  $\langle 100 \rangle$  axis. Then, we tilted the sample at  $6^\circ$  and subsequently a complete scanning on the  $z$  axis was performed in order to identify the  $\{100\}$  and  $\{110\}$  planes. Following, a channeling spectrum was acquired at  $15^\circ$  with respect to the  $\{100\}$  plane. In the sequence, a random spectrum was recorded. This procedure was repeated at  $30^\circ$ ,  $60^\circ$  and  $75^\circ$  with respect to the  $\{100\}$  plane by selection of a new irradiation area on the sample. Typical backscattering spectra are shown in Fig. 1 for  $H^+$  projectiles. Finally, the whole procedure was repeated with the  $H_2^+$  and  $H_3^+$  molecules taking care that the energy per amu and the current per particle were the same in all experiments.

According to the procedure outlined in [10], the electronic stopping power can be determined directly from the spectrum of backscattered particles by fitting a particular function that takes into account the contribution of the dechanneling particles. The results of the  $H^+$ ,  $H_2^+$  and  $H_3^+$  channeling stopping powers as a function of the incident angle are displayed in Fig. 2. Each datapoint was obtained by the following procedure: in the first place, we have done the average between the four individual azimuth measurements. Then, we have taken the mean value between each pair of symmetrical tilt angles (e.g.  $+0.2^\circ$  and  $-0.2^\circ$ ). This mean value is the one plotted at each side of the

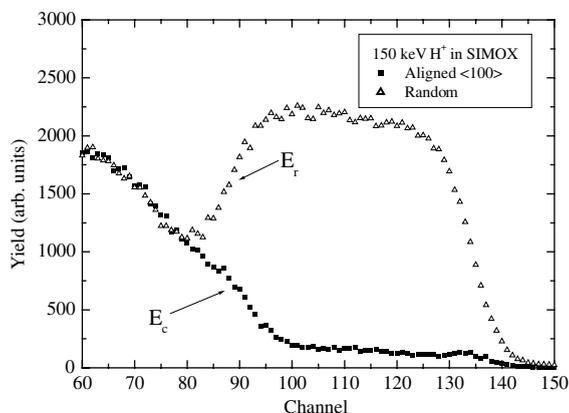


Fig. 1. RBS/channeling spectra of the SIMOX target. Open triangles stand for random spectrum and the squares stand for the  $\langle 100 \rangle$  channeling spectrum.

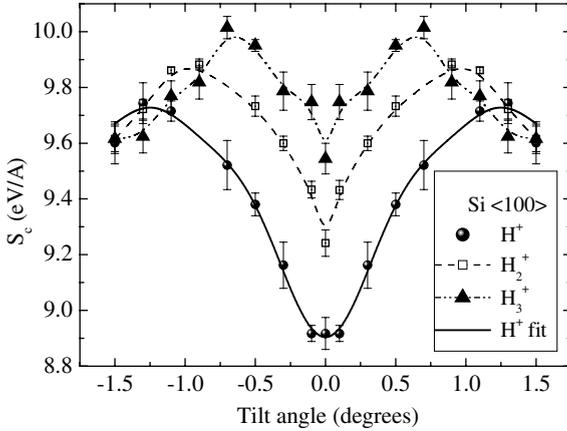


Fig. 2. Channeling stopping power as a function of the tilt angle. Full circles correspond to  $H^+$  ions, open squares correspond to  $H_2^+$  molecule and full triangles to  $H_3^+$  molecules. The full line is the fitting to the  $H^+$  data (using two gaussian functions) and the other lines are plotted only to guide the eyes.

curve shown in Fig. 2. The error bars represent the statistical uncertainty involved in the overall procedure described previously. As can be observed, the  $H^+$  scanning shows the most pronounced dip and the most distant shoulders with respect to the center of the channel, followed by the one corresponding to  $H_2^+$  and, finally, the shallower one belonging to the  $H_3^+$  molecule, having the nearest shoulders.

### 3. Data analysis and results

For the analysis of the  $H_2^+$  and  $H_3^+$  stopping powers we assume that the molecular stopping power  $S_m$  as a function of the tilt angle  $\Psi$  and averaged over the azimuthal angles has three components,

$$S_m(H_n^+, \Psi) = S(H^+, \Psi) + \Delta S_{vic}(H_n^+) + \Delta S_{exp}(H_n^+, \Psi) \quad (1)$$

with  $S(H^+, \Psi)$  corresponding to uncorrelated  $H^+$  fragments, the contributions due to the vicinage and the Coulomb explosion effects  $\Delta S_{vic}$  and  $\Delta S_{exp}$ , respectively. The sum of these last two contributions is shown in Fig. 3. For this sake, we have subtracted from the  $H_2^+$  and  $H_3^+$  distributions the one corresponding to the  $S(H^+, \Psi)$ . Here we

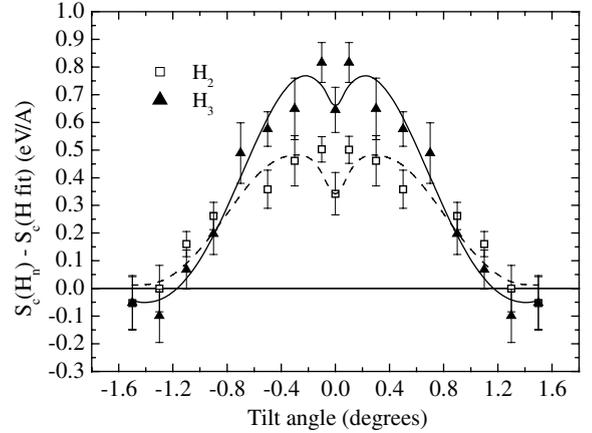


Fig. 3. Molecular  $H_2^+$  and  $H_3^+$  stopping powers after subtracting the  $H^+$  contribution.

have used a fitting curve for  $S(H^+, \Psi)$  (displayed in Fig. 2 with a full line).

The contributions of the vicinage and Coulomb explosion effects on the molecular stopping power shown in Fig. 3 have a peculiar shape as a function of the tilt angle and can be easily interpreted by invoking the angular compensation rule proposed by Lindhard [11]. The angular average over all tilt and azimuthal angles washes out any crystalline effects and its mean value shall correspond to the one for an amorphous target. In that way, if there is an enhancement of the energy loss for  $\Psi \approx 0$  due to vicinage and/or the transversal heating during the molecular breakup, there should be a compensating effect that shall lead to a decreasing of the energy loss at some other tilt angle  $\Psi$ . Physically, this decreasing at larger tilt angles can be explained in terms of a rechanneling of one of the  $H^+$  fragments. In the case of  $H_2^+$  projectiles at 150 keV/amu, the aperture angle (after the explosion) is about the angle where the shoulders appear.

The angular shape in Fig. 3 can be modeled by considering the following conditions:

- $\Delta S_{vic}$  should be nearly independent of  $\Psi$ . It has been shown in [4] that the vicinage effect has a weak dependence on the projectile orientation. Therefore, it basically affects the whole  $S_m(\Psi)$  angular distribution without modifying its shape;

- (b) the mean value of  $\Delta S_{\text{exp}}$  has to be equal to zero in accordance to the Lindhard's angular compensation rule [11] ( $\int d\Psi \sin(\Psi) \times \Delta S_{\text{exp}}(\Psi) = 0$ );
- (c) for larger  $\Psi$ , the function  $\Delta S_{\text{exp}}$  should vanish. In non-aligned conditions, the ion flux distribution is nearly uniform. Therefore, the effect of the Coulomb explosion should be negligible;
- (d)  $\Delta S_{\text{exp}}(\text{H}_n^+, \Psi)$  should be an even function.

We found that the simplest function (for  $\Psi \ll 10^\circ$ ) that fulfills all these conditions is

$$\Delta S_{\text{exp}}(\text{H}_n^+, \Psi) = C \sin(k |\Psi| + \gamma) e^{-\frac{k|\Psi|(1-\cos\gamma)}{\sin\gamma}}, \quad (2)$$

where  $C$ ,  $k$  and  $\gamma$  are constants. These constants and  $\Delta S_{\text{vic}}$  (tilt independent) have been used as fitting parameters.

In Fig. 3, with full and dashed lines the fitting of Eq. (2) to the  $\text{H}_2^+$  and  $\text{H}_3^+$  results are shown. As can be observed our proposed ansatz Eq. (2) describes very well the experimental data. The contribution of the Coulomb explosion for  $\Psi = 0$  are following ones: for  $\text{H}_2^+$ ,  $0.18 \pm 0.05 \text{ eV/\AA}$ ; and for  $\text{H}_3^+$ ,  $0.49 \pm 0.05 \text{ eV/\AA}$ . These values correspond to about 2% and 5% of the stopping of  $\text{H}^+$  in Si.

The maximum contribution of  $\Delta S_{\text{exp}}$  to the stopping power occurs at about the critical angle  $\Psi_c$ . It reaches  $0.3 \text{ eV/\AA}$  for  $\text{H}_2$  and  $0.6 \text{ eV/\AA}$  for  $\text{H}_3$ . Finally, as already mentioned, the negative values of  $\Delta S_{\text{exp}}$  indicate a rechanneling of one of the individual ions. The Coulomb explosion, acting as a astigmatic lens for the ion beam [6], can be the responsible for the rechanneling effect. In fact, it can redirect part of the molecule fragments that otherwise would not be channeled if the beam had no angular dispersion.

In summary, by performing a simple RBS/channeling experiment with  $\text{H}^+$ ,  $\text{H}_2^+$  and  $\text{H}_3^+$  projectiles on a SIMOX target we were able to evaluate the Coulomb explosion contribution to the total stopping power. With this aim, we have performed, for each angle, each azimuth, and each projectile, stopping power measurements. The experiments

were repeated at four different azimuths and an average has given the  $\text{H}^+$ ,  $\text{H}_2^+$  and  $\text{H}_3^+$  final stopping powers as a function of the incident angle. Subsequently, we have subtracted from the  $\text{H}_2^+$  and  $\text{H}_3^+$  stopping powers the  $\text{H}^+$  contribution. In a next step, we have fitted the obtained results with an expression that contains the vicinage and the Coulomb explosion effects. By assuming that the vicinage effect does not affect significantly the shape of the resulting distributions and the angular compensation rule of Lindhard, we were able to extract the  $\text{H}_2^+$  and  $\text{H}_3^+$  Coulomb explosion contributions to the total stopping power.

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